

[54] **CHEMICAL PROCESS YIELDING
STIMULATING EMISSION OF VISIBLE
RADIATION VIA FAST NEAR RESONANT
ENERGY TRANSFER**

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Related U.S. Application Data

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[51] Int. Cl.⁵ **H01S 3/095**

[52] U.S. Cl. **372/89; 372/39;
372/92; 372/56**

[58] Field of Search 372/89, 1, 39, 92, 50

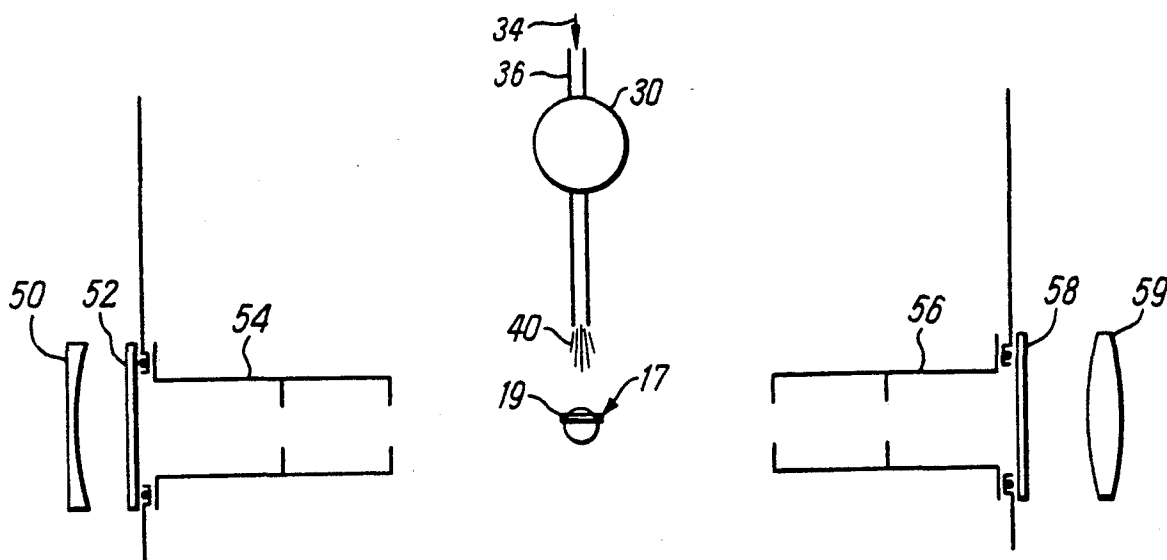
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[57] **ABSTRACT**

A laser oscillator at 535 nm produced by a chemical process yielding stimulated emission of visible radiation via fast near resonant intermolecular energy transfer comprising the steps of reacting a reaction with a first metal vapor to form metastable states of metal oxides or metal halides, transferring the energy from the metastable states to receptor metal atoms by means of near resonant energy transfer to form electronically excited receptor metal atoms in an inverted configuration, and, through multiple reflection, allowing for the repeated passage of light through the inverted gain medium so as to produce oscillation.

21 Claims, 5 Drawing Sheets



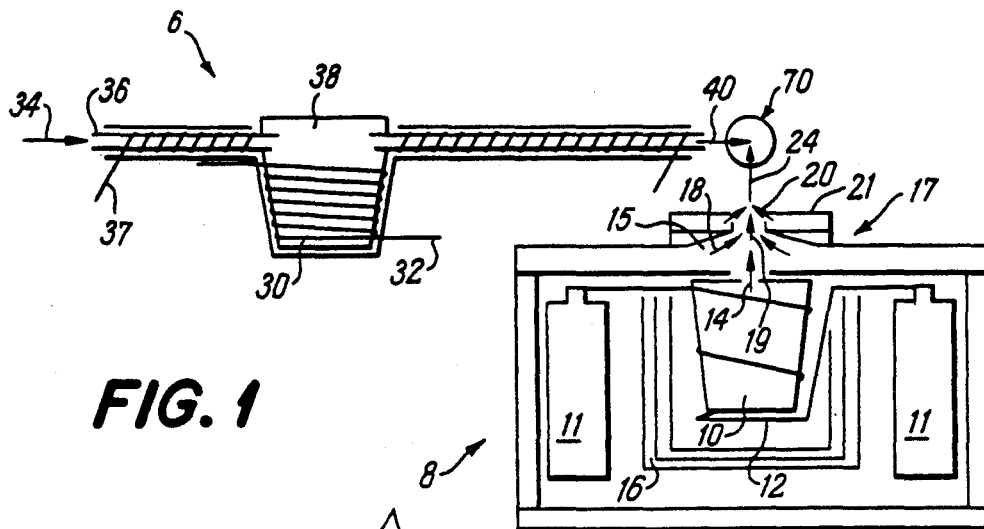


FIG. 1

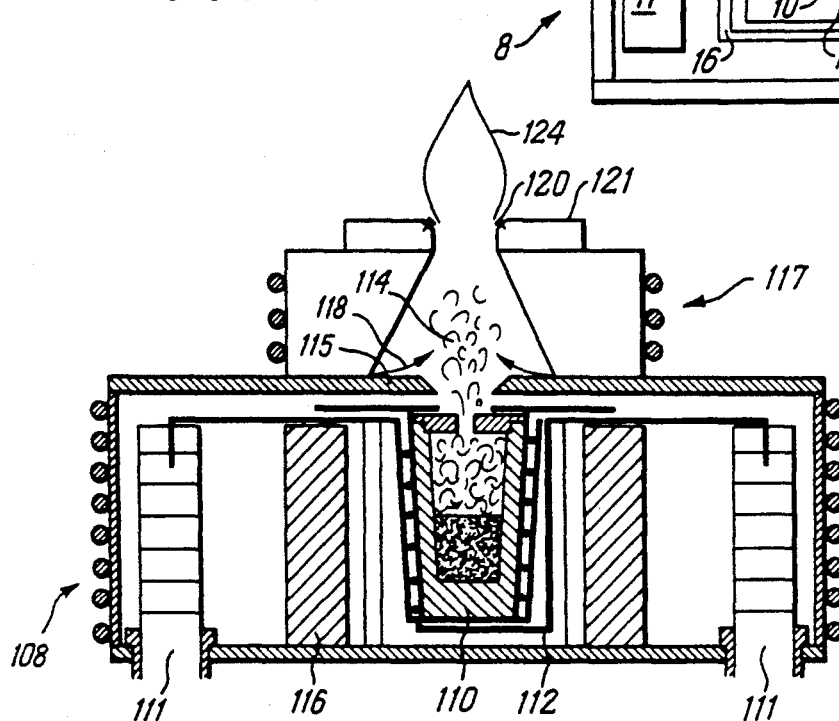


FIG. 2

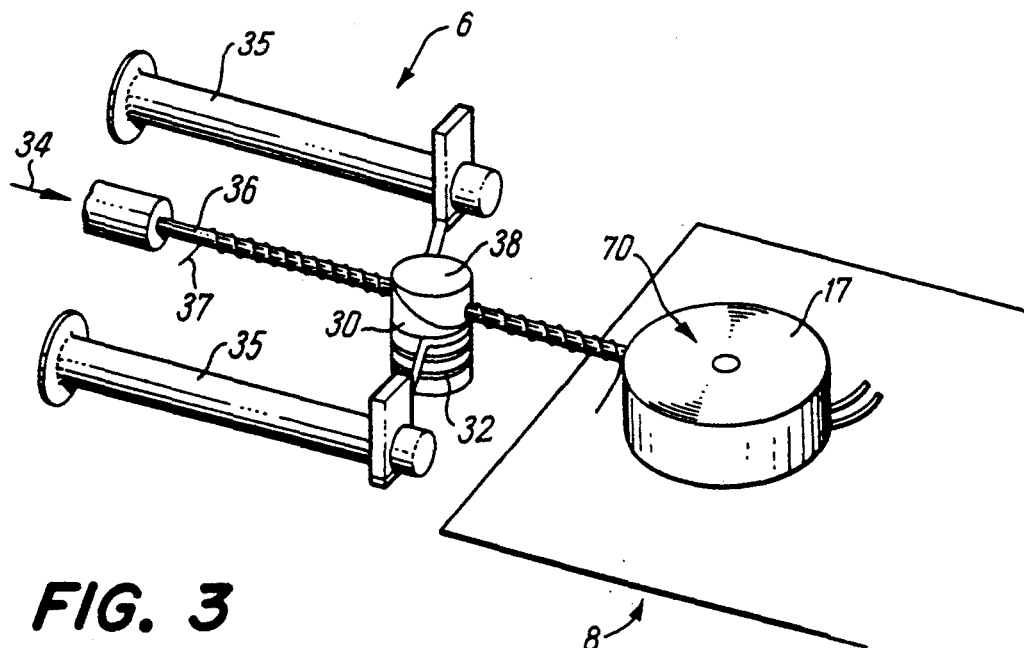


FIG. 3

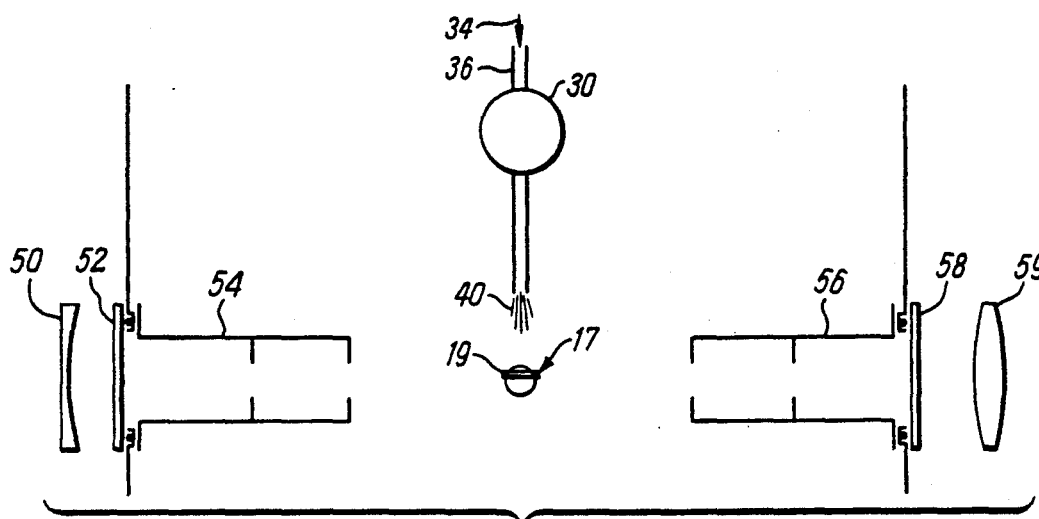


FIG. 4

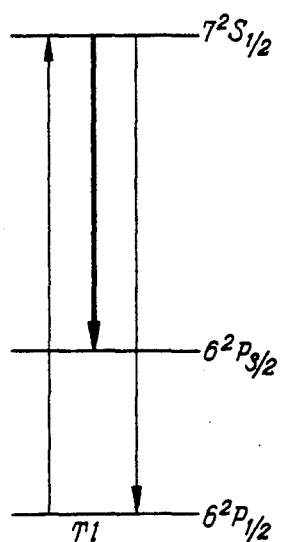


FIG. 5a

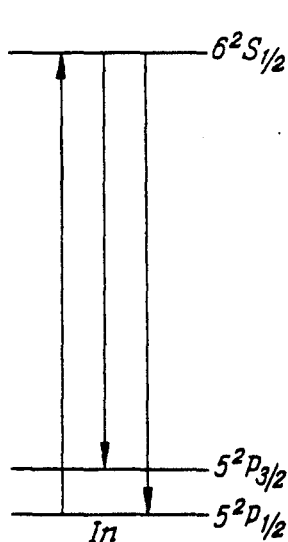


FIG. 5b

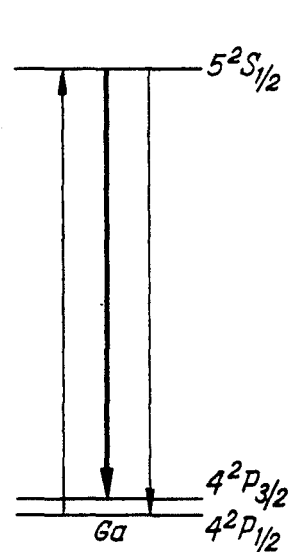


FIG. 5c

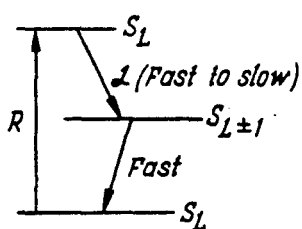


FIG. 6a

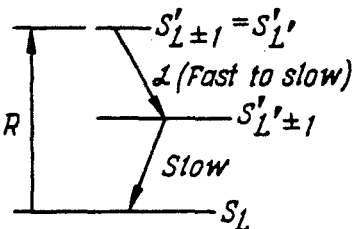


FIG. 6b

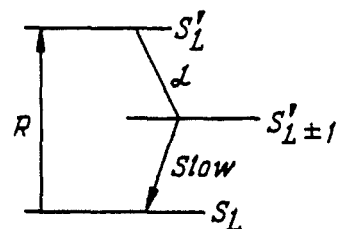


FIG. 6c

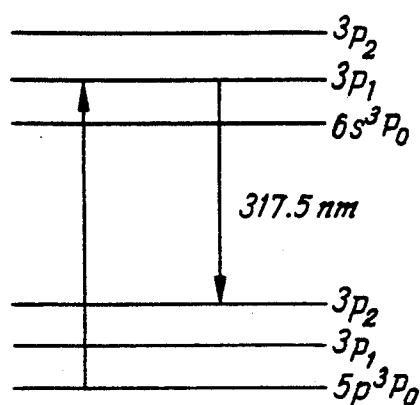


FIG. 7a

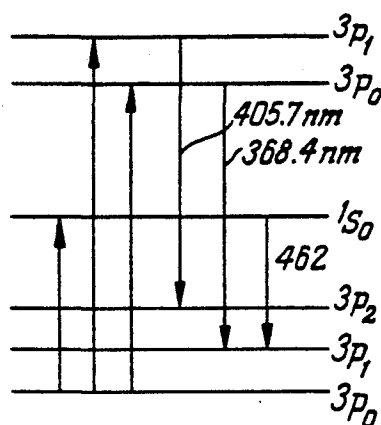


FIG. 7b

Chemiluminescent spectrum
 $\text{Sc} + \text{F}_2 \rightarrow \text{ScF}^* + \text{F}$

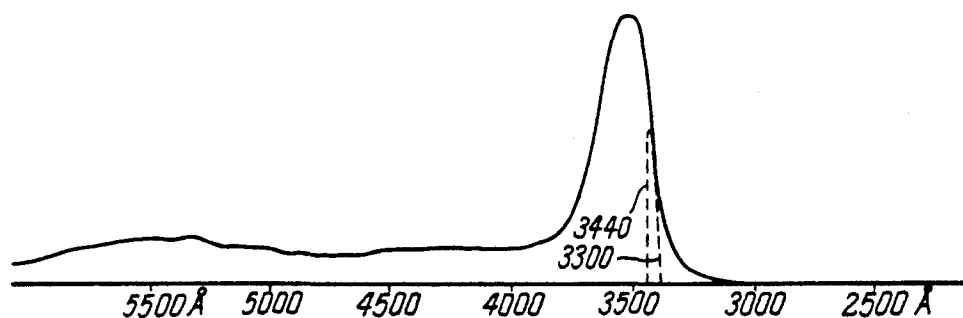


FIG. 8a

Chemiluminescent spectrum
 $\text{Y} + \text{Br}_2 \rightarrow \text{YBr}^* + \text{Br}$

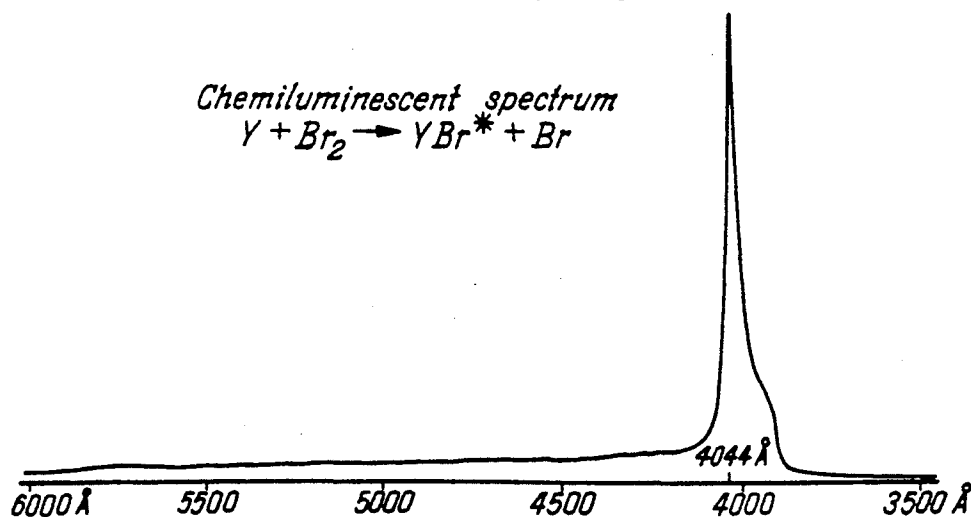


FIG. 8b

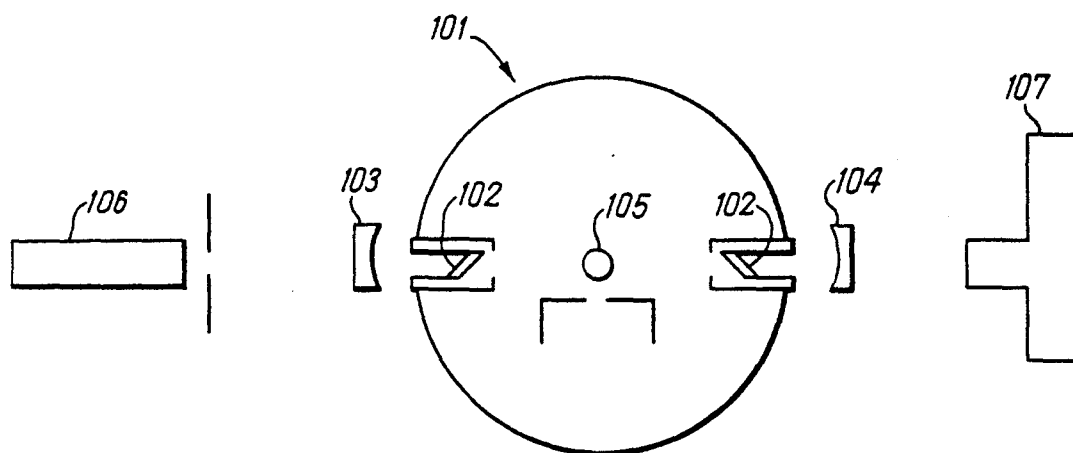


FIG. 9

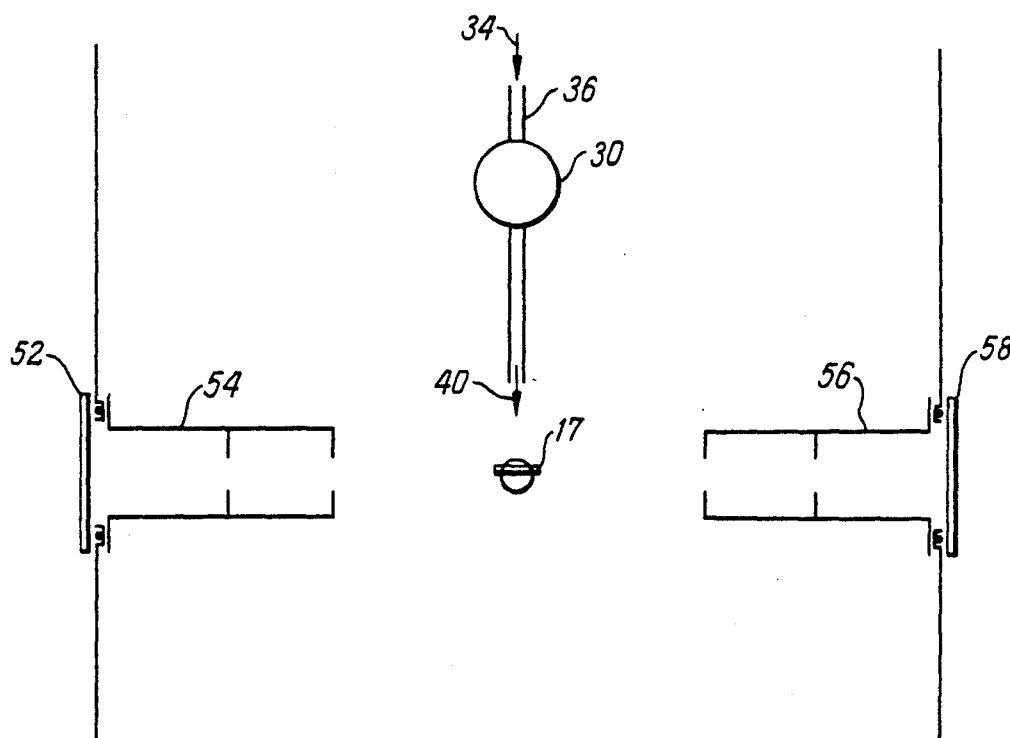


FIG. 10

CHEMICAL PROCESS YIELDING STIMULATING EMISSION OF VISIBLE RADIATION VIA FAST NEAR RESONANT ENERGY TRANSFER

STATEMENT OF GOVERNMENT RIGHTS

The invention was made with Government support under Contract No. DAAL03-88-K0154 awarded by U.S. Army Research office. The Government has certain rights in the invention.

This application is a continuation-in-part of application Ser. No. 07/296,512, filed Jan. 12, 1989.

BACKGROUND OF THE INVENTION

The development of purely chemically driven visible lasers had been an unattained goal for at least the past twenty years. Visible chemical lasers are difficult to implement because the chemical processes which directly establish requisite population inversions leading to lasing are rare. Until now, the only chemical processes known and used to establish population inversions have lead to lasing primarily in the infrared spectral region, the furthest extension toward the visible being to 1.3 μ . The infrared chemical laser systems cannot be used to produce visible lasing, and an entirely new set of processes had to be developed to create the purely chemically driven visible laser of the present invention. As a result, the process of this invention is wholly not described in any prior art and, as such, there is no known relevant prior art.

Classical laser operation in general requires a population inversion in which the upper energy level associated with the lasing transition is more populated than is the lower energy level on which the transition terminates. Laser oscillation can be established in an optical cavity which allows photons to be reflected back and forth and interact with each other so as to build up the intensity of the radiation. A select group of lasers, including primarily N_2 which operates in a pulsed mode under amplified spontaneous emission, does not require such an optical cavity as the photon amplification is so large that sufficient intensity is produced without the necessity of mirrors. A further technique involves stimulated Raman pumping in which an intense laser beam is converted into a beam of another frequency by coherent Raman stimulation in a two or more step scattering process. Each of these lasers operates on electrical energy.

Various types of infrared chemical lasers are known of which there are two main types. The first type involves the mixing of an oxidizer and a fuel gas to produce a continuous output. The mixture is activated by an electrical discharge or by thermal decomposition induced using arc heaters or combustors. The activated mixture produces a reaction initiating species, the reaction sequence eventually leading to a population inversion and lasing involving one of the constituents of the mixture. The second type uses premixed fuels and oxidizers which are activated by flash photolysis with an electron beam or a pulse discharge. However, once the mixture is ignited, it may present flameout and detonation problems and may be difficult to extinguish. Hydrogen halides and carbon monoxide are the two main types of molecules used as lasing species in these chemical lasers.

A typical chemical laser is disclosed in U.S. Pat. No. 4,553,243. This laser operates by expanding the reactant gas mixture continuously through a supersonic nozzle

and applying a pulsed electrical discharge to initiate the chemical reaction resulting in the production of the lasing species. The frequency of the electrical pulses can be adjusted so as to regulate the frequency of the laser.

The gas mixture is introduced on the fly, usually at pressures from a few Torr to multi-atmospheres. This laser does not operate as a purely chemically driven system as it requires an electrical discharge to initiate the chemical reaction. More importantly, this laser only operates in the infrared region and cannot produce visible lasing.

A chemical oxygen-iodine laser using iodine chloride as a reactant gas is disclosed in U.S. Pat. No. 4,563,062. In this device, iodine chloride is vaporized and entrained in argon gas. This gas mixture is directed into a laser cavity where it is mixed with singlet oxygen. Upon mixing with singlet oxygen, the iodine chloride dissociates into atomic iodine and atomic chlorine. Subsequently, the atomic iodine is excited to a lasing state through collisions with the singlet oxygen. These lasers are typically operated with a laser cavity pressure in the range of 1-3 Torr. Using the oxygen-iodine system one produces a laser which operates at 1.3 μ , in the infrared region. Such systems have not been developed to produce visible lasers.

A typical semiconductor laser, based on the gallium arsenide semiconductor, is disclosed in U.S. Pat. No. 4,446,557. This type of laser requires that an external electric field to be applied using an electrode located on the semiconductor layers. When the external electric field is applied, photons are created which resonate among the semiconductor energy levels so as to produce lasing action. The laser produced by a semiconductor system is at a much longer wavelength than the laser of the present invention and operates among much lower lying energy levels.

Self-pulsing, semiconductor lasers which have a pulsed, rather than a continuous wave output beam are also disclosed in U.S. Pat. No. 4,446,557. However, it is difficult to reliably reproduce devices having a very high pulse repetition rate with extremely good temporal stability. Certain geometries for the semiconductor cavity length directly related to the principal noise resonance wave length have been suggested to alleviate this problem. Self-pulsed semiconductor lasers have the same limitations as the typical semiconductor lasers mentioned above.

Metals having sufficient vapor pressures at relatively low temperatures can be made to lase. To create the vapor pressure necessary for lasing action, metals have been heated in electric or gas fired furnaces to approximately 1675°-1875° K. The large amounts of metal vapor required to make such a laser practical require considerable electric power for heating, thus making the resultant laser very bulky and not readily susceptible to mobilization. The use of a gas fired furnace, which is more mobile than an electric furnace, lessens this problem to some degree but the system is still bulky. The use of either an external oven or discharge heating to produce the high temperatures of between 1675° and 1875° K. makes it difficult to construct the fast discharge circuitry needed for excitation of other self-terminating neutral metal laser transitions. Using metal halides helps to reduce the temperature requirements to some degree. The invention which we report for certain applications requires oven systems operating at temperatures in excess of 1600 K., however, for example, the

source of silicon or germanium atoms required to produce metastable storage states which make operative several of the systems of the present invention may be obtained from gaseous silane or germane oxidation reactions as noted in the following Detailed Description of the Invention. A number of the metal atoms required as energy recipients and subsequent atomic lasants can be obtained from sources operating at temperatures considerably less than 1600 K. In addition, the technology needed for operation at the higher temperatures required to operate the particular sources considered here is readily available.

Another commonly used method for creating metal vapor lasers is to sputter the metal atoms from a cathode of the desired material. Control of the sputtering process has been achieved by entraining the sputtered metal in a gas stream so as to create a metal vapor beam. This metal vapor beam is directed into an optical cavity where a separate electrical discharge system excites the metal vapor. Generally, the metal vapor beam is passed through a ring shaped electrode in order to minimize the electrical discharge necessary to excite the lasing constituents in the beam. Metal vapor lasers are not premised on chemical processes such as those reactions used in the present invention. First, there is no chemical reaction. Second, when based on the metal halides, they generally employ dissociation processes caused by an external laser. Third, they are largely operative in the infrared region with only a few examples operative at shorter wavelengths.

Metal halide pulsed lasers capable of simultaneously providing a plurality of output beams oscillating at discrete wavelengths in the visible and near infrared portions of the spectrum are disclosed in U.S. Pat. No. 4,607,371. Such a plurality of output beams is obtained through the dissociative excitation of a number of vaporized metal halides composed of the Group II B metals. Excitation is achieved either by photo-dissociation or by dissociation through collisions with energetic electrons produced in a transverse discharge or by an electron beam generator. The power of such lasers can be enhanced by using isotopically pure metal halide salts rather than their naturally abundant counterparts. As such, this laser relies on a dissociation process caused by an external laser and not a chemical reaction.

Chemically driven visible lasers offer attractive alternatives to their infrared counterparts; however, the development of a chemically pumped system lasing in the visible region, while occupying the interest of researchers for almost two decades, represents a difficult problem whose solution has met with little success until the present invention. This invention focuses on the development, extension, and detailed quantification of visible chemical laser systems and the demonstration and quantification of laser amplification and oscillation across the visible and ultraviolet regions employing purely or primarily chemically pumped systems. The development of such devices necessitates innovative approaches to the generation of electro-magnetic radiation. To implement these approaches, we take advantage of the unique features associated with certain high cross-section, highly selective exothermic reactions and several new insights gained in the study of ultra fast energy transfer processes in small high temperature molecules.

Efforts toward the goal of a visible chemical laser oscillator are to be encouraged for not only is there reason to pursue these systems for their potential high

gain, but also, given similar power levels, a device based on an electronic transition holds advantage over those based on infrared transitions in that the size of the device may be smaller, the power consumption efficiency larger, and the optics considerably simplified relative to the infrared. Devices once constructed and optimized can play an important role in a diversity of field based operations. Finally, chemical lasers once developed are inherently more efficient than systems based exclusively on electrical power input.

SUMMARY OF THE INVENTION

Until the present invention was developed, stimulated emission of visible radiation had not been produced through a solely chemical pumping process. The process described in the present invention requires no optical or electrical pumping to create a population inversion leading to a subsequent superfluorescent event. The present invention is a visible chemical laser system based on an efficient near resonant (greater than or equal to gas kinetic) energy transfer involving metastable excited states of metal or semimetal monoxides, formed in the reaction of metal or semimetal atoms with ozone, N_2O , or NO_2 , and Group IIIA $2P_{1/2}$ atoms in the lower spin orbit component of their ground electronic states (X_{gl}). This energy transfer populates the Group IIIA $2S_{1/2}$ excited state (X^*) creating a population inversion which subsequently provides the basis for a superfluorescent event. The concept of fast near resonant energy transfer to subsequently lasing metal or semimetal atoms is readily extended to several other near resonant processes described in detail below which we also claim as a subset of this invention.

The present invention is the first known to create a population inversion in a final subsequently lasing constituent via a near resonant intermolecular energy transfer from another constituent formed in a select chemical reaction. The inversion is manifest in the stimulated emission of visual radiation.

The superfluorescent laser pulse created on the thallium $72S_{1/2}-62P_{3/2}$ transition whose duration in a self-terminating system is $\sim 5ns$ FWHM. This single pass superfluorescent system is converted to a multipass laser oscillator (3% output coupling) with a corresponding increase in laser output power correlating with a substantial increase in the ratio of superfluorescence to fluorescence and the display of significant directionality.

DESCRIPTION OF THE DRAWINGS

FIG. 1 is an elevation, partly in section, of a typical reaction configuration associated with the present invention including an oven assembly utilized in forming MO metastable molecules;

FIG. 2 is an elevation, partly in section, of an alternate embodiment of an oven assembly associated with the Present invention utilized in forming MO metastable molecules;

FIG. 3 is a schematic view of the oven source used to produce a vapor of gallium, indium or thallium atoms in near-resonant energy transfer experiments;

FIG. 4 is a top view of the oven configuration and photon path to the spectrometer in the near-resonant energy transfer system;

FIG. 5 are representative energy level diagrams for thallium, SiO a-X, and GeO a-X systems of the present invention;

FIG. 6 is the energy level diagram for three level laser systems created using the fast intermolecular energy transfer concept of the present invention as a pump (R) to create an inversion and lasing (L);

FIG. 7 are the energy levels and pump and lasing transitions for the tin and lead receptor systems of the present invention;

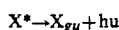
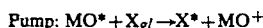
FIG. 8 is the chemiluminescent emission from the $\text{Sc} + \text{F}_2$ and $\text{Y} + \text{Br}_2$ reactions depicting the selective formation of metastable metal halide excited states.

FIG. 9 is a top view of the reaction zone configuration in the oscillation system;

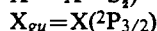
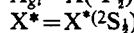
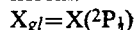
FIG. 10 is a top view of the oven configuration and photon path in the oscillation system.

DETAILED DESCRIPTION OF THE INVENTION

The present invention employs the concept demonstrated by the superfluorescence associated with an atomic transition of a Group IIIA metal or semimetal atom. Although Group IIIA includes semimetals, such as boron, the term metal will be used throughout this specification to denote all of the elements, whether metals or semimetals, in Group IIIA. The metal atom electronically excited to a level X^* , which has been chemically pumped through energy transfer from low-lying metastable states of select metal oxides or metal halides, undergoes a superfluorescent transition to X_{gu} (upper spin orbit component of the Group IIIA atom ground electronic state [or a low-lying excited electronic state]). The following near resonant intermolecular transfer scheme is appropriate to the Group IIIA atoms and exemplary of the present invention:

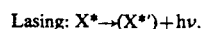


Wherein:



In this scheme, the metal oxide (MO) excitation is transferred to a Group IIIA atom (X), specifically and preferably thallium, but also including gallium, leading to a pumping from the ground state X_{gl} to energetically accessible electronically excited levels, X^* , of the Group IIIA atom. FIG. 5 exemplifies energy level diagrams appropriate to the thallium atom, as it interacts with SiO (a-X) and GeO (a-X) levels. These diagrams indicate the energy match-up between the various a-X energy increments associated with $v'=0$ SiO and GeO $a^3\Sigma^{30}$, and the thallium electronic level structure. The electronically excited thallium atoms, X^* , pumped by energy transfer subsequently undergo transitions from X^* to X_{gu} .

The potential laser transition X^* to X_{gu} can be made superfluorescent. For several systems (following) we may extrapolate on the Group IIIA atom concept replacing the ground state spin orbit components ($X^2P_{1/2}$) and $X(^2P_{3/2})$ with ground (X_g) and low-lying electronically excited (X^*) metal atoms, the superfluorescent transition corresponding to:



The development of electronic transition chemical lasers generally requires a two-step approach in which chemical energy is provided and stored in the first step and then this chemical energy is transferred to an appro-

priate laser medium in the second step. In the present invention, the second step occurs in a fast, near-resonant, energy transfer from the metastable states of metal oxides or metal halides preferably SiO and GeO (although silicon and germanium generally are considered to be semimetals, the term metal will be used throughout this specification to denote both silicon and germanium as well as all of the elements, whether metals or semimetals, in Group IVA), to readily lasing atoms which include the Group IIIA metals, preferably thallium but also including gallium. This transfer is facilitated primarily by the near resonance of the metastable metal oxide energy levels and the energy required to pump the electronically excited state energy levels X^* , for example, of the thallium and gallium atoms as illustrated in FIG. 5. The energy matchup between various a-X energy increments associated with $v'=0$ and the thallium atomic energy level structure in FIG. 5 are indicated (Δ) in wavenumbers. In the case of thallium, the strong ground state spin orbit coupling provides a large energy increment between the $X^2P_{1/2}$ (X_{gl} at 0.000 cm^{-1}) and $X^2P_{3/2}$ (X_{gu} at 7793 cm^{-1}) spin orbit levels greatly facilitating the simple pumping scheme outlined above.

Given that the metastable concentration of metal oxides and receptor metal atoms required to produce a threshold electronically excited state, X^* , concentration through efficient energy transfer is produced, the upper level spin orbit component X_{gu} acts as the terminating level in a three level laser scheme.

The metal to be oxidized to form metastable excited states is heated to a temperature producing a vapor pressure between approximately 10^{-1} and 2 Torr. For silicon the operating temperature is preferably between about 1800 to 2000 K.; for germanium the temperature range is approximately 1600–1850 K. These temperatures produce a sufficient concentration of metal vapor for the energy transfer-lasing process especially after entrainment in a carrier gas. The metal vapor 14 is entrained in an inert gas flow 18. This metal vapor/inert gas mixture 24 is introduced into the reaction zone 70. One method of accomplishing this is diagrammed in FIGS. 1–4.

The metal to be oxidized is held in a graphite crucible 10 which is heated by a tungsten basket resistive heater 12 insulated by extensive tantalum and zirconia baffling 16. Power to the heater 12 is supplied through electrodes 11. The crucible 10 is generally brought to temperature over a two-hour period. The metal vapor 14 effusing from the crucible 10 is entrained in an argon flow 18 being supplied through a circular directed slit configuration 15 below the lower portion of a concentric ring injector 17, and forming a combined metal vapor/inert carrier gas flow 19.

The metal to be oxidized can also be oxidized in the oven assembly 108 of FIG. 2 which is an alternate, but equally appropriate, embodiment of the oven assembly 8 of FIG. 1. Oven assembly 108 incorporates parallel features to oven assembly 8, and is likewise numbered in parallel. For example, crucible 10 in FIG. 1 is replaced with crucible 110 in FIG. 2, electrode in FIG. 1 is replaced with electrode 111 in FIG. 2, and so on. This oven assembly 108 operates in a similar manner to oven assembly 8, the major difference being the precise orientation of the entrainment configuration 117 including ring injector 21 which are modified to optimize the gas flow and entrainment in a given pumping configuration.

The effective vapor pressure of the metal vapor in the argon flow will be a little higher than the vapor pressure mentioned above. The crucible used need only be a low porosity, non-reactive container in which a Group IVA (as previously noted, all of the elements in Group IVA, whether metal or semimetal, will be denoted in this specification as a metal) or other appropriate metal may be held during heating. As a practical matter, the temperature to which the metal which reacts to form the metastable energy storage state is heated and the volume of argon flow are adjusted to achieve the desired result.

Referring specifically to FIG. 3, the metal whose atomic vapor forms the atomic receptor, X_{g1} , and subsequent lasing is typically placed in a second crucible 30 and heated to a temperature producing approximately a 10^{-1} to 10 Torr vapor pressure. For the Group IIIA metal, thallium, one generally wishes to operate the particular system described in FIGS. 1-4 at a temperature close to 1100 K. This operating temperature provides approximately a 1 Torr vapor pressure or higher, which is the preferred concentration of the atomic receptor metal vapor, neglecting the effects of a cooler entraining inert gas stream. An inert carrier gas 34 is passed through the top of the crucible 38, forming a flow of an atomic (ex: Group IIIA metals thallium and gallium) metal vapor/inert gas mixture 40 which intersects the metal vapor/inert gas flow 24, in the reaction zone 70. The maximum vapor pressure of atomic receptor metal vapor is just below the concentration at which the oxidation reaction which forms the metastable metal oxide (ex: rxn of Group IVA metals Si or Ge) may become self-quenching. Any inert gas may be used as carrier gas 34, however, argon is preferred because of its molecular weight and lower cost. Further, any containers may be substituted for the crucibles so long as they are non-reactive with the metal to be oxidized or the atomic receptor metal and will withstand the temperatures required for operation. A significant variety of crucible materials are available for this purpose, their use being dictated primarily by cost.

As an example, thallium can be placed in an aluminum oxide crucible 30 which is heated to a temperature of 1100 K. using a tantalum wire resistive heating configuration 32, operating at about 12 amperes at 10 volts alternating current. When using gallium or indium as the Group IIIA metal, higher temperatures are needed. This is now accomplished by supplying up to 70 amperes at 2.5 volts a/c to a tungsten wire basket heater 32 through electrodes 35. Further the alumina tubes 36 bringing inert carrier gas to the crucible are still heated using a tantalum wire configuration 37 of the design of FIG. 1 is used. This same modification can also be used for thallium. The crucible 30 generally was brought to temperature over a two-hour period.

Again using thallium as an example, the argon gas 34 is passed through alumina connecting tubes 36 and through the top 38 of the crucible 30. The thallium vapor/argon mixture 40 forms a flow which is directed perpendicular to and intersects with the metal vapor/inert gas flow 24, described above in the reaction zone 70. The alumina tubes 36 and top of the crucible 38 are heated by tantalum wire resistive heating 37 to eliminate the possibility that a cooler argon entrainment flow 40 might result in severe condensation problems. The thallium atoms have a velocity in excess of 7×10^4 cm-sec⁻¹ and travel across the reaction zone 70 in less than approximately 14 ms, producing a Group IIIA metal

atom flux in the reaction zone 70 ranging from 10^{17} (10^{-2} Torr) to 5×10^{20} cm²-sec for the totality of measurements conducted thus far with the thallium system. An extremely efficient energy transfer from the MO* metastable states to the X_{g1} atoms excites the X_{g1} atoms to the X^* state. As an alternative to the perpendicular flow described above, the two metal vapor flows may be brought to the reaction zone in parallel through a concentric injector comprised of one tube located inside a second tube. The size of the reaction zone depends upon a number of factors including the relative locations of the two metal source configurations and the pumping of the metal/vapor inert carrier gas mixture for the two flows 19 and 40.

At this point, the system is capable of either lasing or fluorescing. In order for the system to lase, the receptor metal vapor/inert gas flow 40 must first mix with the metal vapor/inert gas flow 24 in the reaction zone 70. When this mixing has occurred, an oxidant 20, which for metastable metal oxide formation is typically NO₂, N₂O, or O₃ dictated by the particular metal reaction and energy transfer partners is introduced into the reaction zone through the upper portion 21 of the concentric ring injector system 17. For example, for the Ge-GeO-Tl system the oxidant of choice is O₃; for the Si-SiO-Pb system the oxidant of choice is NO₂. The introduction of the oxidant 20 triggers the lasing sequence as it produces the oxidation process.

FIG. 4 represents a top view of the oven configurations and optical system for measuring a superfluorescent laser pulse. It includes two light baffles (54, 56) to shield the detectors from oven system blackbody radiation, two high quality quartz windows (52, 58) and a focusing lens 59. In order to increase the ratio of superfluorescent to fluorescent light intensity a "half oscillator" configuration employing the 100% reflective mirror 50 has been employed. This represents the first stage in the creation of a full oscillator multipass configuration.

When using ozone, one method of controlling the O₃ flow to the reaction zone 70 is to freeze the O₃ down on a silica gel trap at dry ice temperature, then passing a flow of O₃ through a short stainless steel tube to the concentric ring injector 17 using a triggered pulsed valve or a manual needle valve (not shown). If the O₃ concentration obtained is not high enough, the O₃ flow can be backed with an inert carrier gas, such as helium, in order to increase the O₃ vapor pressure in the reaction zone 70. In this manner, for example, the vapor pressure of O₃ can be readily adjusted to the requisite pressure, if desired, up to about 1 Torr.

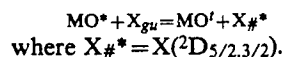
Preferably, the vapor pressure of oxidant in the reaction zone will be less than 1 Torr and preferably in the range of about 5×10^{-2} to 2×10^{-1} Torr. The oxidant flow may be either continuous or pulsed. However, once the reaction - energy transfer sequence leading to superfluorescent emission from the receptor atom produces a loss of population inversion due to a filling of the low laser level (ex: upper spin orbit component or low-lying electronic state in thallium or gallium) lasing ceases for a time period dictated by the diffusion of potentially metastable atoms populating the terminal laser level. If the terminating laser level is not metastable, lasing action can be maintained for a time frame considerably longer than the radiative lifetime associated with the $X^* \rightarrow X_{gu}$ ($X^* \rightarrow X^*$) transition with which the emission of laser light is associated.

The key to effecting stimulated emission of visible radiation using the present invention is the formation of a sufficient concentration of metastable or triplet excited state metal oxides or, as described below, metal halides by chemical reaction. These metastable states must be sufficiently long-lived to first store energy from the chemical reaction and then must be capable of transferring that energy on a sufficiently fast time scale to a medium capable of emitting at least a portion of this energy as visible radiation. Several metal oxidations lead to formation of metastable storage states of the Group IVA metal oxides, such as GeO and SiO. These metastable states have sufficiently long-lasting radiative lifetimes so as to meet the above criteria, the radiative lifetime of the GeO and SiO $a^3\Sigma^+$ and $b^3\pi$ states being approximately 10^{-2} to 10^{-1} seconds. We have further found that the ground states of the Group IIIA metals, thallium and gallium, are readily pumped through fast near resonant energy transfer from the metastable states of such Group IVA metal oxides. The ratio of the Group IIIA metal, such as thallium, to the Group IVA metal oxide, such as GeO*, is preferably between 1:1 and 1.5:1.

The mixing sequence is important in these systems. To induce lasing action, the metal vapor which is oxidized to initiate the superfluorescent sequence must first be mixed with the receptor atom which will produce the stimulated emission event. The oxidant then triggers the lasing sequence as it is introduced into the reaction zone in a pulsed or continuous fashion. If this sequence is not followed, it is more likely that the system will fluoresce and not lase. In order for the system to display energy transfer primarily in the form of receptor atom fluorescence as opposed to lasing, the metal vapor-argon flow 24 is mixed with various oxidants 20 as it is introduced into the reaction zone 70. As above, using the Ge-GeO-Tl system as an example, the oxidant 20 mixes with the metal vapor/inert gas flow 24 via a concentric ring injector system 17. As the metal vapor/inert gas flow 24 mixes with the oxidant as it enters the reaction zone 70, a metal oxide flame 24 is produced. More specifically, an entrained Group IVA metal flux 14 under multiple collision conditions reacts with the oxidant 20 to produce MO*. If at this time the Group IIIA metal vapor/inert gas flow 40 is introduced into the metal oxide flame, the probability for observing a superfluorescent, as opposed to fluorescent event, is considerably diminished as a result of mixing requirements and the population of the terminating laser level. For example, an entrained Si flux under multiple collision conditions reacts with an oxidant, specifically NO₂, N₂O, or O₃, to produce SiO metastables or an entrained Ge flux under multiple collision conditions reacts with an oxidant, specifically N₂O or O₃, to produce GeO metastables. For most efficient usage, these metastables must usually be formed after the two metal flows employed are mixed previously in order to insure lasing action.

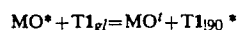
At the concentrations of receptor atoms used in the present invention exemplified by thallium or gallium, self-absorption by ground state (X_{gl}) atoms can occur. For example, the relative intensities of the fluorescence features resulting from the X^* to X_{gu} and X^* to X_{gl} transitions in thallium are approximately in the ratio of 8 to 1 at the highest Group IIIA metal concentrations. This ratio should be 1.16 to 1 if saturation effects are not present. Indeed this ratio is obtained at a thallium vapor pressure close to 10^{-2} Torr in the present system. The

overwhelming determinant for the observed ratio at higher thallium metal flux is ground state self-absorption. Additionally, other metallic emission features may also be observed. The thallium system is characterized by emission features resulting at least in part from the chemical pumping of efficiently produced X_{gu} atoms formed via the initial pump - superfluorescence, fluorescence sequence described above followed by:



We exemplify the detection of superfluorescence and hence lasing action using the Ge-GeO-Tl system. To detect superfluorescent emission from thallium at 5351 Å, the emission signal from the appropriate Tl transitions was detected in time using a digital oscilloscope. The Ge-O₃ metatheses, chosen to form the metastable metal oxide and the subsequent energy transfer are controlled initially with the O₃ flow through a triggered pulsed valve or manual needle valve (not shown). The observed fluorescence signal from the thallium 5351 Å line was measured at approximately 50 mV while the time-resolved superfluorescence was measured at approximately 500 mV (a lower bound due to response time) with a FWHM of approximately 5 ns, in excess of ten times the normal $^2S_{1/2} - ^2P_{3/2}$ fluorescent intensity. However, no superfluorescence signal spike occurs from the Tl fluorescence features at 3777 Å and 3525 Å or with any excited metal oxide emission features.

The Tl_{gu} atoms are created in the reaction-energy transfer zone, as is indicated by the thallium emission at 3525 Å resulting from the transition from Tl₁₉₀* to Tl_{gu}. The Tl_#* energy levels can be populated through pumping of Tl_{gu} in a near resonant energy transfer from MO*, the energy increment also being close to that for pumping from Tl_{gl} to Tl*. The intensity, pressure, and reactant concentration dependence of the emission at 3525 Å are consistent at least in part with the presence of a Tl_{gu} concentration in the reaction zone, which subsequently is pumped to the Tl_#* manifold through the reaction:



Although we have exemplified the present invention using primarily thallium as the lasing element, another Group IIIA element, gallium, may also be used as a lasing species chemically pumped through energy transfer from a Group IVA metal oxide, such as SiO and GeO. The P level spin-orbit splitting of the ground electronic state of the gallium atom is considerably smaller than that of thallium ($G_a = 826 \text{ cm}^{-1}$) thus allowing for increased thermal population of X_{gl} . This complicates the lasing process but the reactivity of gallium with ozone and especially the relative reactivity of gallium X_{gu} atoms with ozone versus that reaction involving X_{gl} atoms aids the lasing process. In addition, cold, nonreactive CO, CO₂, or N₂ can be used as an entrainment gas so as to relax these X_{gu} atoms before energy is transferred to them from the excited metal oxides. A direct extension of the thallium and gallium laser systems corresponds to the Si-SiO-Pb system where the oxidant reacting with Si to produce the SiO metastables is NO₂.

Fast Near Resonant Intermolecular Energy Transfer

The fast near resonant intermolecular energy transfer concept which we have outlined, has focused on a

pump of thallium and gallium excited states, X^* , but can be readily extended if we consider that all of these systems should operate on the basis of energy storage created by a reaction whose exothermicity exceeds 2.5 eV (or approximately 58 Kcal/mole). The near resonant energy transfer concept can be used as an efficient pump for somewhat longer lived emitters (allowing a further increase in energy storage in a laser cavity) Replacing the $X^2P_{3/2}$ spin orbit component with a low-lying electronic state (X^*) in a variety of atomic receptors, we generalize the near resonant intermolecular energy transfer pump-amplification concept to a broader range of lasing configurations. Further, it is feasible to replace the metastable states of SiO and GeO with selectively formed and long-lived states of the Group IIIB halides.

As applied to thallium, gallium and indium atoms, the results we have obtained indicate the importance of the near matchup of SiO and GeO energy level increments with the energy level spacing of the receptor atoms of interest and (2) the potential for operation over a rather wide range of intermediate ($P_{3/2}$) and ground state (P_1) level separations. With the extension to low-lying electronic states, the three level excitation scheme offers a considerable range of radiative configurations (lifetimes) commensurate with a variety of potential lasing transitions.

There are a number of three level schemes where, in large part, an energy transfer pump from the metastable states of SiO and GeO can be accomplished with a near resonant intermolecular energy transfer exciting a state only weakly optically coupled to the lowest ground state level of a receptor atom. Lasing may occur to an intermediate state whose optical coupling to the pumped state varies from strong to only moderate; the radiative lifetimes associated with these transitions vary from several nanoseconds to several microseconds, the transitions within themselves covering a range of visible wavelengths. Further, within this pumping scheme, lasing occurs to a state which is either weakly or strongly coupled to the ground state, this combination providing a range of pump-amplifier schemes. Such systems offer pump transitions which are closer in resonance to SiO and GeO than are the energy levels of the Group IIIA atoms, Ga, Tl, and In. Thus, they are potentially more efficient. To be specific, the schemes indicated in FIG. 6 are deemed direct extrapolations of the present invention. The intermediate level is unoccupied before the energy transfer pump is applied.

With reference to FIG. 6(a) (employing in place of either thallium or gallium as atomic receptors, alkali and alkaline earth metals, Sc, V, Fe, Ni, or Y) both the total spin and orbital angular momentum are maintained in the energy transfer-pump transition so that these levels are not strongly optically coupled. To first order both transitions between the upper and intermediate and intermediate and ground levels are strongly allowed. However, for certain atomic receptors, the transitions from the upper to intermediate level are found to be parity violating, providing a sharp increase in the radiative lifetime associated with the laser transition. Further, we anticipate intermediate cases where the nature of the atomic "wavefunctions" for those levels associated with the lasing transition affects the atomic transition of interest but not to the extent of a parity violation. With this framework, we predict radiative lifetimes associated with the laser transition ranging from 10^{-8} to 10–5 seconds allowing one to draw a desired balance

between increased mixing and energy storage and the coupling of the amplifying medium to the laser cavity.

Referring to FIG. 6(b), employing as receptor atoms Ti, Cr, Ba, and to FIG. 6(c) employing as receptor atoms Sc, Ti, V, Ni, Sr, and Ba, the energy transfer pump transition involves a violation of the spin angular momentum selection rule, $\Delta S=0$, as one excites an upper level of spin S' different than that of the ground level (S). To first order, transitions between the upper and intermediate level are strongly allowed whereas transitions from the intermediate to ground state level are spin forbidden. Thus, the intermediate level would be metastable, and these two systems will operate in a pulsed mode analogous to the thallium and gallium atom systems which we have detailed. However, again due to the nature of those "wavefunctions" associated with the upper and intermediate atomic levels, the radiative lifetimes associated with those transitions from the upper to the intermediate level would vary across a wide range. In each of the schemes illustrated by FIG. 6 we employ fast near resonant energy transfer from Group IVA metal oxides, such as GeO and SiO.

Finally, we note a fourth pump-amplification configuration in direct analogy to the thallium and gallium systems. Particularly, we focus on the pumping of several strongly or weakly optically coupled transitions primarily in lead, but also in tin, where we again make use of the large spin orbit splittings (FIG. 7) characterizing the ground electronic 3P states of these atoms. We employ near resonant energy transfer from SiO and GeO metastable states to pump the 3P_0 component for the ground electronic states of the lead and tin atom viz.

Again, the pump transitions indicated are in closer resonance to SiO and GeO metastable level separations than are the energy levels of the Group IIIA atoms Ga, Tl, and In. Thus, the energy transfer pump will be inherently considerably more efficient.

Generation of High Si and Ge (SiH, GeH) Concentrations from Silanes and Germanes

There are several approaches which can be used to generate substantial concentrations of the metastable states of SiO and GeO from the oxidation of gas phase silane, germane, disilane, and digermane. We can generate both the metal monohydrides (SiH, GeH) and metal atoms (Si, Ge). SiH and GeH can react with O_3 to yield the metastable states of the metal monoxide. We have produced the requisite metastable states via the reaction of these four hydrides with N_2O and O_3 in a purely chemical albeit complex reaction sequence. Chemiluminescence from the low-lying electronic states of SiF, SiH, and SiD has been observed by several workers studying the reactions of SiH_4 , Si_2H_6 , and SiD_4 with molecular fluorine. Some of these experiments have involved the use of close to 1 millimole/sec of reactants in a lab based operation. From these studies it is apparent that one can readily generate 1.25×10^{17} /sec (SiH+Si) direct chemiluminescent precursors for a consumption of 2.5×10^{20} SiH_4 molecules in a fluorine limited configuration. This, in turn suggests the potential production of SiO and GeO metastables in high yield from the corresponding reactions with N_2O , O_3 , and possibly NO_2 .

While reactive stripping represents one route to the desired SiO molecules, this approach is in general oxidant limited. An alternative primarily chemically based hybrid approach combining and extending the techniques which have proved useful in silicon vapor depo-

sition with the advantages of reactive encounters may be used as a means of obtaining substantial SiO and GeO metastable state concentrations. Substantial silicon atom and subhydride concentrations have been generated employing pyrolytic techniques and low pressure plasmas as a means of generating hydrogenated amorphous silicon films. Further, simple hollow cathode techniques have been used recently to generate metal atom concentrations well in excess of $10^{14}/\text{cm}^3$.

Sufficient concentrations of the silanes and germanes can be broken down in a high pressure plasma or electric discharge to the atoms and metal monohydrides (possibly in combination with pyrolysis) and subsequently oxidized to produce the metastable states of SiO and GeO. The technology already exists to considerably extend the established successful lower pressure generation schemes. While microwave plasmas can now be operated effectively in the 10 to 20 Torr pressure regime, electric discharge techniques have now been perfected at very much higher pressures and hence reactant flows, sustaining an electric discharge of 10–20 kV/cm-atm in pulsed operation. These field levels which should be sufficient to provide for the efficient stripping of the silanes and germanes to SiH, GeH, Si, and Ge at high pressure can be improved to even higher operational levels.

This combinatorial approach which can be used to generate substantial concentrations of SiO and GeO metastables, at least two orders of magnitude in excess of those already obtained, is by no means a slight variant on electric discharge techniques. In fact, it represents a discharge enhanced primarily chemical process. This approach allows one to bring the SiH radicals (GeH), silicon atoms (Ge) and receptor atoms into intimate contact for the purpose of premixing before the oxidation-energy transfer pump sequence is initiated.

Further Extensions

A further extension of the intermolecular energy transfer concept of the present invention focuses on the halogenation reactions of the Group IIIB metals Sc, Y, and La. The grid in Table 1 presents twenty examples of Group IIIB-halogen reactions. As we exemplify in FIG. 8, we have found extremely selective production of long-lived halide excited electronic states in many of the scandium and yttrium reactions which are easily observed because of an extremely high quantum yield for excited state formation. Analysis of the temperature dependence for several representative reactions indicates that selective excited state formation proceeds by a direct mechanism with negligible activation energy.

The data in FIG. 8, indicates the range of wavelengths over which the selectively formed excited electronic state emits. The ScF selective emission feature peaks at 3500 Å whereas the corresponding YBr feature peaks at 4040 Å. Thus, for the combination of scandium and yttrium oxidation with the homonuclear and heteronuclear halogen molecules, we cover the range 3400–4150 Å. For those metastable states based on yttrium halogenation, this range is well suited for intermolecular energy transfer to subsequently lasing gallium, indium, or thallium atoms. Incorporating both the scandium and yttrium halogenations, we further increase and extend the range of possible laser systems within a wide group of potential receptor atoms.

TABLE I

Group IIIB - Halogen Reactions Studied-Product Formation			
	Sc	Y	La
F ₂	ScF ^s	YF ^s	LaF
Cl ₂	ScCl ^s , ScCl ₂ ^c	YCl ^s	LaCl
Br ₂	ScBr ^s , ScBr ₂ ^c	YBr ^s	—
I ₂	ScI ^s , ScI ₂ ^c	YI ^s	—
SF ₆	ScF ^s	YF ^s	LaF
ClF	ScF ^s > ScCl	YCl ^s > YF ^s	Laf, LaCl
ICl	ScCl ^s > ScI	YCl ^s > YI ^s	—
IBr	ScBr ^s > ScI	YBr ^s > YI ^s	—

^s = selective monohalide emission

^c = continuous emission due to dihalide.

Extension of the Thallium Atom Laser Amplifier Concept to Full Oscillation

To achieve superfluorescent lasing action and subsequent oscillation: (1) copious quantities of the SiO and GeO metastable energy storage states in "high quantum yield" chemiluminescent reactions must be produced; (2) sufficient thallium atom $X^2P_{3/2}$ concentrations and simultaneously minimal $X^2P_{3/2}$ concentrations must be produced before the pump and lasing processes above are made to occur; and (3) near resonant (highly efficient) intermolecular energy transfers between the SiO and GeO metastable states and those levels accessible via excitation of the lowest $^2P_{1/2}$ spin orbit component of the thallium atom must occur.

Given that the requisite concentrations are obtained (for the Ge-O₃-Tl system, the concentrations in the reaction zone are $\text{Ge} \geq 2 \times 10^{14} \text{ atom/cm}^3$, $\text{O}_3 \geq 7 \times 10^{14} \text{ /cc}$, $\text{Tl } X^2P_{1/2} \geq 10^{16} \text{ atom/cc}$), the upper $^2P_{3/2}$ component of the ground state thallium atom can act as the terminal level in a three level laser provided that this state is not populated before the lasing pump-sequence is initiated. Thallium is used, as discussed above, because its $X^2P_{3/2}$ level is 7793 cm^{-1} above the $X^2P_{1/2}$ level thus preventing thermal excitation of this level ($N_{3/2} = 10^{-10}(N_1)$ at $T = 1100 \text{ K.}$) during vaporization, subsequent entrainment, and passage to the reaction zone, and insuring its population exclusively through the reaction sequence. Thallium also will not compete with those reactive processes forming the metastable states of SiO and GeO. The TlO bond is weak, if not unstable, with any oxidation processes considered for its production being highly endothermic.

The single pass laser amplifier systems described above can be converted to a multipass laser oscillator configuration in which the light associated with the stimulated emission process is made to pass several times through the gain medium (reaction-energy transfer zone) producing a corresponding increase in the laser output power and the ratio of superfluorescence (or super-radiance) to fluorescence in the display of a significant directionality.

In order to convert superfluorescent (vs. amplified spontaneous emission) laser amplification to full multipass laser oscillation in which we obtain an enhanced lasing output relative to the single-pass amplifier configuration, a full cavity configuration 101 shown in FIG. 9 consisting of two cavity mirrors 103, 104, one near totally reflective mirror 103 and one with between 0.1 and 3% transmissivity at the wave length of interest dictated by the particular reaction-energy transfer configuration in use 104. In the case of the high gain Ge-GeO-Tl system whose amplification and oscillation are obtained from Tl atoms, these mirrors are about 100

and about 97 to about 99% reflective (about 1 to about 3% transmittance) at 535 nm; however for the Si-SiO-Na amplifier system, based upon amplification on a sodium atom transition with a much lower gain coefficient, the required mirrors are about 100% and about 99% to 99.8% reflective at 565 nm.

Within the constraints of the concave mirror locations are placed Brewster angle windows 102 (for example, Mellos Griot BK-7 ($n=1.519$) on axis on opposite sides of the reaction-energy transfer-amplification zone and oriented at a specific angle to minimize reflectivity losses in the cavity for that wavelength associated with the transition producing laser action. For example, the angle for the Ge-GeO-T1 system minimizes reflectivity losses at 535 nm, whereas that for the Si-SiO-Na system minimizes reflectivity losses at 565 nm. These Brewster angle windows may be replaced equivalently by anti-reflector coated windows perpendicular to the cavity axis and operative in the wave length ranges for the transition of the amplifying medium (e.g., CVI laser optics).

The concave mirrors 103, 104 are chosen to operate in a stable cavity configuration such that $0 < g_1 g_2 < 1$ where $g_1 = 1 - L/r_1$ and $g_2 = 1 - L/r_2$, L being the separation between the mirrors (cavity length) and r_1 and r_2 being the radii of curvature of each mirror, respectively. For example, in the T1 based system, two concave mirrors of 2 m radius of curvature, and a separation between the mirrors of 30 cm satisfy the criteria for a stable laser cavity. These 2 m radius of curvature mirrors might be chosen primarily because of their ready availability in the laboratory (Nd-YAG laser mirrors) and the requirement to create a stable laser cavity configuration. The 30 cm cavity length is dictated by the particular apparatus construction of FIG. 4. As the cavity length is varied over a range where $L < r_1$ or r_2 , the stable resonator configuration is maintained; however, the number of possible passes through the gain medium, dictated primarily by the radiative lifetime associated with the transition and the excited stage, X^* , pumping rate for the atomic receptor atom from which lasing is obtained, is affected. For example, light traveling at about 3×10^{10} cm/sec traverses a 30 cm cavity in 1 nanosecond (10^{-9} seconds). This allows for 3 to 5 amplified passes through the gain medium in the Ge-GeO-T1 system of FIGS. 4, 9, and 10 before lasing action ceases. The number of passes would increase by a factor of three for a 10 cm cavity. For the Si-SiO-Na system at 568 nm, a factor of 10 increase in the radiative lifetime allows for an increase to at least 10 to 40 passes through a similarly defined amplifying gain medium. Several mirror combinations which create a stable resonator configuration are possible, especially those which most effectively match the laser mode volume and the gain medium volume. The cavity dimensions are by no means rigid and the mirror configurations and output coupling can be modified to obtain optimization. Thus, the mirror separations, radii of curvature, location of the gain medium relative to the cavity mirrors and the degree (percent) of output coupling can be modified to optimize several possible systems; including all amplifier configurations considered previously.

For the thallium system, we have used a cavity length of about 30 cm and a reaction (amplifying medium) zone 105 between about 1-1.5 cm long as dictated by the particular experimental configuration of FIGS. 4, 9 and 0. The reaction zone and cavity dimensions are

operational, but by no means optimal and can be improved with a number of cavity 101 configuration modifications including a lengthening of the amplifier zone and those others mentioned above. A HeNe laser 106 which is used to align the cavity also specifies the region of the much shorter amplification zone 105 which will be sampled. The lever arm formed by the HeNe laser beam about the amplification zone 105 central to the current 30 cm cavity 101 has a significant moment as this system by no means constitutes an optimal cavity configuration.

The T1 laser system has also been probed temporally using a 125 MHz digital oscilloscope in conjunction with two monitoring configurations. In the first, the laser cavity 101 is aligned with the incident slit of a monochromator 107 set at 535.1 nm for the T1 laser system, operating at 1 nm resolution, and placed ~ 20 cm behind the output coupling mirror 104 of the laser cavity. The output from an RCA 4840 photomultiplier used previously to monitor superfluorescence is again sent to a digital oscilloscope used previously to monitor the T1 atom amplifier system. The signal recorded in full cavity configuration (FIG. 9) is approximately a factor of ten more intense than the superfluorescent laser amplifier pulse recorded in the absence of a full laser cavity and focusing lens (FIG. 10). Further, the superfluorescence to fluorescence ratio increases to well over 100.

In a second experimental monitoring configuration for the T1 system, the monochromator was replaced with a fast photodiode (LeCroiz Model 40D Optical Detector—risetime < 1.0 ns) (not shown) placed 20 cm from the laser cavity 101. The output from this photodiode was again sent directly to a 125 MHz digital oscilloscope. The signal level recorded for the T1 system was approximately three times that from a 1 mW HeNe laser impinging directly onto the photodiode and placed 2 cm from it.

In both of the experiments described, when the laser cavity mirrors 103, 104 are walked from their aligned positions, oscillatory lasing action for the T1 system at 535 nm can be made to cease and only amplification is detected. The combination of results obtained in the present study demonstrates the temporal and directional characteristics which are to be associated with a multipass visible chemical laser oscillator which we now claim. To distinguish, we can modify a single pass amplifier system (FIG. 10) which for T1 can be made to create a multiple pass laser oscillator through the introduction of cavity mirrors aligned in a stable cavity configuration on opposite ends of a Brewster angle windowed laser cavity configuration set for 535 nm (FIG. 9).

It now becomes apparent that the above described chemical processes yielding stimulated emission of visible radiation via fast near resonant energy transfer and the apparatus for carrying said process are capable of obtaining the above-stated objects and advantages. It is obvious that those skilled in the art may make modifications to the process or the apparatus or both, without departing from the spirit of the invention or the scope of the subjoined claims.

We claim:

1. A chemical process yielding stimulated emission of visible radiation via fast rear resonant intermolecular energy transfer comprising the steps of:

(a) providing a first source of vapor selected from the group consisting of metal or semimetal vapor to be

- reacted with an appropriate reactant to produce a metastable excited state reaction product which serves as an energy storage medium;
- (b) providing a second source of atomic vapor selected from the group consisting of metal or semi-metal vapor to serve as receptor atoms to receive the energy from said metastable excited state of the reaction product;
- (c) providing a source of reactant to react with the first source of vapor in a highly exothermic reaction which liberates energy exceeding 2.5 eV; said reactant being selected from the group consisting of ozone, nitrogen oxide, nitrogen dioxide, and halides;
- (d) chemically reacting the reactant and the first source of vapor to form metastable states of a final reaction product;
- (e) transferring energy stored in the metastable states of the said reaction product to the second source of atomic vapor serving as receptor atoms by means of near resonant energy transfer to form electronically excited receptor atoms in a population inversion relative to a lower level of excitation of said receptor atoms; and
- (f) providing a mirror configuration having first and second mirrors and windows associated with said mirrors, wherein said mirrors are chosen such that
- $$0 < g_1 g_2 < 1,$$
- where
- $$g_1 = 1 - L/r_1$$
- $$g_2 = 1 - L/r_2$$
- L = separation between the mirrors (cavity length)
- r_1 = radius of curvature of said first mirror
- r_2 = radius of curvature of said second mirror
- $L < r_1$ and
- $L < r_2$,
- said first and second mirrors being placed on opposite ends of an enclosed amplification zone such that the inverted population and multiple reflections allows for oscillation via the repeated passage of light through the inverted gain medium operative at the wavelength of the amplified transition associated with the population inversion created among the energy levels of the receptor atoms in a stable cavity configuration,
- said windows being operative in the wavelength ranges for the transition of the amplifying medium defined by the length of the cavity and radii of curvature of said mirrors, and
- said mirror configuration causing light associated with the amplified transition to undergo multiple passes through the inverted gain medium yielding oscillation of stimulated emission of visible radiation.
2. A process as defined in claim 1 wherein the receptor atoms are non-reactive with the selected reactant.

3. A process as defined in claim 1 wherein the receptor atoms are selected from the group consisting of atoms capable of receiving gas kinetic or greater near resonant energy transfer from the metastable states of the oxide or halide.
4. A process as defined in claim 1 wherein the receptor atoms are selected from the group consisting of thallium, gallium, lead, alkali, and alkaline earth metals, scandium, vanadium, iron, nickel, yttrium, titanium, chromium, barium and strontium.
5. The process as claimed in claim 4 wherein said receptor atom is thallium and $L < 30 \text{ cm} < r_1, r_2$.
6. The process as claimed in claim 4 wherein said receptor atom is sodium and $L < 30 \text{ cm} < r_1, r_2$.
7. A process as defined in claim 1 wherein said first source of vapor is selected from the group consisting of germanium and silicon.
8. A process as defined in claim 1 wherein said first source of vapor is silicon and the reactant is selected from the group consisting of nitrogen oxide and nitrogen dioxide.
9. A process as defined in claim 1 wherein said first source of vapor is silicon and the reactant is nitrogen dioxide.
10. A process as defined in claim 1 wherein said first source of vapor is germanium and the reactant is ozone.
11. A process as defined in claim 1 wherein said first source of vapor is selected from the group consisting of scandium, yttrium, and lanthanum and the reactant is a halogen.
12. A process as defined in claim 1 wherein step (a) includes first providing a source of a first element selected from the group consisting of metals and semimetals and heating said first element to form the first source of vapor to be reacted.
13. A process as defined in claim 12, wherein said first element is heated to a temperature producing a vapor pressure of approximately 10^{-1} to 2 Torr.
14. A process as defined in claim 1 wherein step (b) includes first providing a source of a second element selected from the group consisting of metals and semimetals and heating said second element to form said second source of vapor to serve as receptor atoms.
15. A process as defined in claim 11, wherein said second element is heated to a temperature producing approximately a 10^{-1} to 10 Torr vapor pressure.
16. A process as defined in claim 1, wherein said mirror configuration comprises two mirrors.
17. A process as defined in claim 16, wherein said mirrors are concave inwardly towards each other.
18. A process as defined in claim 12, wherein said mirrors have a radius of curvature of about 2 m.
19. A process as defined in claim 16, wherein one of said mirrors is totally reflective and the other of said mirrors is an output coupler.
20. A process as defined in claim 19, wherein said output coupler has about a 3% transmissivity.
21. A process as defined in claim 1, further comprising a Brewster's windows.

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